## An analytical solution of Lagrange planetary equations valid also for very low eccentricities

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We give an analytical solution of the Lagrange planetary equations expressed in a set of non singular elements for eccentricity: the semi-major axis a, the inclination i, the ascending node  $\Omega$ ,  $\lambda = \omega + M$  where  $\omega$  denotes the argument of perigee and M the mean anomaly, and the elements  $C = e \cos \omega$  and  $S = e \sin \omega$ where e denotes the eccentricity.

The perturbations which can be computed are due to the gravity field of a planet, or to the influence of a third body, through an internal or an external potential. These potentials are expressed with respect to the former set of elements, and in particular use a formulation of the Kaula functions of eccentricity G (Wytrzyszczak, 1986) and H with respect to C and S.

The main study deals with the integration of the differential system verified by C and S, since the metric element e and the angular element  $\omega$  are mixed. The differential system is divided into two parts: a first one which can be analytically integrated in an exact way, the second part being seen as a perturbation of the first one. More precisely, the first part of the differential system verified by C and S corresponds to a harmonic oscillator: our method takes into account, in a first step, long period variations of 4 elements: C,  $\Omega$ , S, and  $\lambda$ .

It is therefore a generalization of the Kaula method(Kaula, 1966) expressed in 'classical' orbital elements, in which only the secular variations of the angles  $\Omega$ ,  $\omega$ , M are computed during a first step.

Moreover, our method enlightens the behavior of the argument of perigee  $\omega$  with time. Like (Cook, 1966), we show that its variations are not really linear with time, in particular for low eccentricities. We therefore use another angle, called  $\theta$ , which is well defined even for eccentricities equal to zero. We use this angle  $\theta$  to integrate the long as well as the short period variations of C and S, and as a consequence of all the other elements.

Here, we show why and how we have divided the differential system verified by  $a, C, i, \Omega, S$  and  $\lambda$  in two components. The first part gives a long period solution  $a_1, C_1, i_1, \Omega_1, S_1$  and  $\lambda_1$ ; the second part gives corrections  $\Delta a, \Delta C, \Delta i, \Delta \Omega, \Delta S$  and  $\Delta \lambda$ . There results:  $a(t) = a_1(t) + \Delta a(t), C(t) = C_1(t) + \Delta C(t), \dots$ 

We also show the way used to express  $\Delta a$ ,  $\Delta C$ ,  $\Delta i$ ,  $\Delta S$ , ... in a synthetic formulation.

Like the method built by Kaula, our method is based on a development in powers of the eccentricity. But, there is no supplementary hypothesis induced by the expression of our theory in non singular elements for eccentricity: it can be used for bodies – artificial satellites or planets – with an eccentricity greater than or equal to zero.

Finally, we compare three methods in different dynamic configurations: this method, expressed in non singular elements for eccentricity (with an integration hypothesis based on  $\theta$ ), the method of Kaula, expressed in classical orbital elements (with an integration hypothesis based on  $\omega$ ), and the numerical integration. Tests show that, compared to the method of Kaula, our method is more accurate as soon as the eccentricity of the studied body is smaller than a few  $10^{-2}$ .

## References

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